
Preprint No. M 17/03

**Coupling of definitizable operators in
Krein spaces**

Vladimir Derkach and Carsten Trunk

2017

URN: urn:nbn:de:gbv:ilm1-2017200208

Impressum:

Hrsg.: Leiter des Instituts für Mathematik
Weimarer Straße 25
98693 Ilmenau
Tel.: +49 3677 69-3621
Fax: +49 3677 69-3270
<http://www.tu-ilmenau.de/math/>

COUPLING OF DEFINITIZABLE OPERATORS IN KREĬN SPACES

VLADIMIR DERKACH AND CARSTEN TRUNK

ABSTRACT. Indefinite Sturm-Liouville operators defined on \mathbb{R} are often considered as a coupling of two semibounded symmetric operators defined on \mathbb{R}^+ and \mathbb{R}^- , respectively. In many situations, those two semibounded symmetric operators have in a special sense good properties like a Hilbert space self-adjoint extension.

In this paper we present an abstract approach to the coupling of two (definitizable) self-adjoint operators. We obtain a characterization for the definitizability and the regularity of the critical points. Finally we study a typical class of indefinite Sturm-Liouville problems on \mathbb{R} .

1. INTRODUCTION

Let \mathcal{K} be a Hilbert space with the inner product (\cdot, \cdot) and let J be a linear operator in \mathcal{K} , such that $J = J^* = J^{-1}$. The space \mathcal{K} endowed with the Hermitian sesquilinear form $[\cdot, \cdot]_{\mathcal{K}} = (J \cdot, \cdot)$ is called a *Kreĭn space* and is denoted by $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$, for details see [4, 10] or Section 2.1 below.

The Hermitian sesquilinear form $[\cdot, \cdot]_{\mathcal{K}}$ induces in an obvious way sign type spectrum for linear operators. E.g., a real isolated eigenvalue is of positive type if all the corresponding eigenvectors are positive with respect to $[\cdot, \cdot]_{\mathcal{K}}$. In the last two decades this notion was frequently used in theoretical physics in connection with \mathcal{PT} -symmetric problems, we mention here only [8, 9, 13, 22, 41] and in the study of \mathcal{PT} -symmetric operators we refer to [2, 12, 43, 44].

A self-adjoint operator A in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ is said to be *definitizable* [38], if its resolvent set $\rho(A)$ is nonempty and there exists a real polynomial p such that $p(A)$ is nonnegative in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. If $\alpha_1 < \alpha_2 < \dots < \alpha_N$ is the set of all real zeros of p then there exists a spectral function $E(\Delta)$ of A , which is defined on all intervals Δ , such that the endpoints of Δ do not belong to the set $\{\alpha_j\}_{j=1}^N$, $E(\Delta)$ takes values in the set of orthogonal projections, commuting with A and $E(\Delta)$ is monotone on each interval (α_j, α_{j+1}) . These intervals are classified in [38] as intervals of positive and negative type and the points α_j which separate intervals of different types are called *critical*. A critical point α is called *regular*, if the operators $E(\Delta)$ are uniformly bounded for all small Δ containing α , otherwise it is called *singular*. The set of critical points of A is denoted by $c(A)$, the set of regular

Date: January 23, 2017.

1991 *Mathematics Subject Classification.* Primary 47B50; Secondary 46C20, 47B25.

Key words and phrases. self-adjoint extension, symmetric operator, Kreĭn space, locally definitizable operator, coupling of operators, boundary triple, Weyl function, regular critical point.

The research of the first author was supported by the Deutsche Forschungsgemeinschaft (DFG) under grant no. TR 903/16-1 and Ministry of Education and Science of Ukraine (projects # 0115U000136, 0115U000556).

(singular) critical points of A is denoted by $c_r(A)$ ($c_s(A)$, respectively). The notion of local definitizability of a self-adjoint operator A in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ was introduced in [24, 25], see Section 3 below.

In the present paper the following problem is studied: The problem of the definitizability of the coupling A of two symmetric operators A_+ and A_- and the regularity of its critical points. Remind the definition of the coupling from [42] adapted to the case of Kreĭn spaces. Let a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ be the orthogonal sum $\mathcal{K} = \mathcal{K}_+[\dot{+}]\mathcal{K}_-$ of $(\mathcal{K}, [\cdot, \cdot])$ of two Kreĭn spaces $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}_+})$ and $(\mathcal{K}_-, [\cdot, \cdot]_{\mathcal{K}_-})$, such that the subspaces

$$\mathcal{D}_+ = \{f \in \mathcal{K}_+ \cap (\text{dom } A) : Af \in \mathcal{K}_+\} \quad \text{and} \quad \mathcal{D}_- = \{f \in \mathcal{K}_- \cap (\text{dom } A) : Af \in \mathcal{K}_-\}$$

are dense in \mathcal{K}_+ and \mathcal{K}_- and the restrictions

$$A_+ = A|_{\mathcal{D}_+} \quad \text{and} \quad A_- = A|_{\mathcal{D}_-}$$

are symmetric operators with defect numbers $(1, 1)$ in the Kreĭn spaces $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}})$ and $(\mathcal{K}_-, [\cdot, \cdot]_{\mathcal{K}})$, respectively. The operator A is called a *coupling* of two symmetric operators A_+ and A_- . The coupling A of two symmetric operators A_+ and A_- is not uniquely defined by the above definition. We will make this definition more precise in Theorem 4.4 by using the boundary triple approach developed in [14, 33, 11, 21]. For differential operators with indefinite weights the coupling method was used in [29], and also in [31, 30, 34] to study the similarity problem and in [7] to study definitizability.

The main result of the paper is Theorem 4.6 where conditions for regularity of the critical point $\infty \in c(A)$ are found under the assumptions that the symmetric operators A_+ and A_- admit definitizable and semibounded extensions $A_{+,0}$ and $A_{-,0}$. The proof is based on the K. Veselić criterion of regularity [46] adapted to the case of definitizable operators in [23]. In the case when A_+ and A_- are Hilbert space symmetric operators similar results were obtained in [34] and [15].

Typically, such problems arise in the study of indefinite Sturm-Liouville operators

$$\ell(f)(t) := \frac{\text{sgn } t}{w(t)} \left(-\frac{d}{dt} \left(\frac{df}{r(t)dt} \right) + q(t)f(t) \right) \quad \text{for a.a. } t \in \mathbb{R}, \quad (1.1)$$

where the coefficients r , q and w are real functions on \mathbb{R} satisfy the conditions

- (C1) $r, q, w \in L^1_{\text{loc}}(\mathbb{R})$ and $r, w > 0$ a.e. on \mathbb{R} ,
- (C2) the expression ℓ is in the limit point case at $-\infty$ and at $+\infty$.
- (C3) minimal differential operators B_{\pm} generated by $\pm \ell$ in $L_w^2(\mathbb{R}_{\pm})$ is semi-bounded from below.

The operator A generated by the differential expression (1.1) in the Kreĭn space is the coupling of two semibounded symmetric operators $A_{\pm} := \pm B_{\pm}$. In Proposition 5.1 it is shown that the operator A is definitizable over a neighborhood of ∞ and conditions (4.18) for $\infty \notin c_s(A)$ are formulated in terms of the m -coefficients of the operators B_{\pm} . In the case $w \equiv 1$ the conditions (4.18) are fulfilled automatically, [15]. This fact was proved earlier by another method in [16].

1.1. Notations and preliminaries. By \mathbb{C}_+ we denote the set of all $z \in \mathbb{C}$ with positive imaginary part and we set $\overline{\mathbb{C}} := \mathbb{C} \cup \{\infty\}$ and $\overline{\mathbb{R}} := \mathbb{R} \cup \{\infty\}$.

A complex function m is called a *Nevanlinna function* if m is holomorphic at least on $\mathbb{C} \setminus \mathbb{R}$ and satisfies the following two conditions

$$m(\overline{z}) = \overline{m(z)} \quad \text{and} \quad \text{Im } m(z) \geq 0, \quad \text{for all } z \in \mathbb{C}_+. \quad (1.2)$$

For facts on Nevanlinna functions we refer to [28] and [20, Chapter II].

All operators in this paper are closed densely defined linear operators. For such an operator T we use the common notation $\rho(T)$, $\text{dom}(T)$, $\text{ran}(T)$ and $\ker(T)$ for the resolvent set, the domain, the range and the null-space, respectively, of T . We define the extended spectrum $\tilde{\sigma}(A)$ of A by $\tilde{\sigma}(A) := \sigma(A)$ if A is bounded and $\tilde{\sigma}(A) := \sigma(A) \cup \{\infty\}$ if A is unbounded and we set $\tilde{\sigma}(A) := \overline{\mathbb{C}} \setminus \tilde{\sigma}(A)$.

2. DEFINITIZABLE OPERATORS IN KREĬN SPACES

2.1. Kreĭn spaces. We recall standard notation and some basic results on Kreĭn spaces. For a complete exposition on the subject (and the proofs of the results below) see the books by Azizov and Iokhvidov [4] and Bognár [10]. A vector space \mathcal{K} with a Hermitian sesquilinear form $[\cdot, \cdot]_{\mathcal{K}}$ is called a *Kreĭn space* if there exists a so-called *fundamental decomposition*

$$\mathcal{K} = \mathcal{K}_+ \dot{+} \mathcal{K}_-,$$

which are orthogonal to each other with respect to $[\cdot, \cdot]_{\mathcal{K}}$ such that $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}})$ and $(\mathcal{K}_-, -[\cdot, \cdot]_{\mathcal{K}})$ are Hilbert spaces. Those two Hilbert spaces induce in a natural way a Hilbert space inner product (\cdot, \cdot) and, hence, a Hilbert space topology on the Kreĭn space \mathcal{K} . Observe that the indefinite metric $[\cdot, \cdot]_{\mathcal{K}}$ and the Hilbert space inner product (\cdot, \cdot) of \mathcal{K} are related by means of a *fundamental symmetry*, i.e. a unitary self-adjoint operator J which satisfies

$$(x, y) = [Jx, y]_{\mathcal{K}} \quad \text{for } x, y \in \mathcal{K}. \quad (2.1)$$

If \mathcal{H} and \mathcal{K} are Kreĭn spaces and $T : \mathcal{H} \rightarrow \mathcal{K}$ a bounded operator, the adjoint operator T^+ of T with respect to the Kreĭn spaces \mathcal{H} and \mathcal{K} is defined by

$$T^+ := J_{\mathcal{H}} T^* J_{\mathcal{K}},$$

where $J_{\mathcal{H}}$ and $J_{\mathcal{K}}$ are the fundamental symmetries associated with \mathcal{H} and \mathcal{K} , respectively; the operator T^+ satisfies $[Tx, y]_{\mathcal{K}} = [x, T^+y]_{\mathcal{K}}$ for all $x \in \mathcal{H}, y \in \mathcal{K}$. If A is a densely defined operator in \mathcal{K} then the *adjoint* A^+ of A with respect to $[\cdot, \cdot]_{\mathcal{K}}$ is defined analogously. In fact, if J is a fundamental symmetry on $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ and (\cdot, \cdot) is the corresponding Hilbert space inner product (2.1), then $A^+ = JA^*J$. The operator A^+ satisfies

$$[Ax, y]_{\mathcal{K}} = [x, A^+y]_{\mathcal{K}} \quad \text{for all } x \in \text{dom}(A), y \in \text{dom}(A^+).$$

In analogy with the definitions in Hilbert spaces, A is *symmetric* in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ if A^+ is an extension of A and A is *self-adjoint* in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ if $A = A^+$.

A densely defined operator A is called *nonnegative* in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ if $[Af, f]_{\mathcal{K}} \geq 0$ for all $f \in \text{dom}(A)$. A nonnegative self-adjoint operator in a Kreĭn space can have an empty resolvent set; a specific example is given in [38, Section 1.2] and [10, Example VII.1.5]. But if a nonnegative self-adjoint operator in a Kreĭn space has also nonempty resolvent set, then it has real spectrum only.

An operator A is called semibounded from below in the Kreĭn spaces $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$, if there exists $\alpha \in \mathbb{R}$ such that

$$[Af, f]_{\mathcal{K}} \geq \alpha[f, f]_{\mathcal{K}}, \quad f \in \text{dom}(A).$$

2.2. Definitizable operators. In this section we recall some facts on definitizable operators in Kreĭn spaces. For an overview we refer to [39], see also [37]. For this purpose it is convenient to introduce in Definition 2.1 below the notion of sign-type spectra, cf. [5, 26, 36, 40].

Let A be a closed operator in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. A point $\lambda_0 \in \mathbb{C}$ is said to belong to the *approximative point spectrum* $\sigma_{ap}(A)$ of A if there exists a sequence (x_n) in $\text{dom}(A)$ with $\|x_n\| = 1$, $n = 1, 2, \dots$, and $\|(A - \lambda_0)x_n\| \rightarrow 0$ if $n \rightarrow \infty$. For a self-adjoint operator A in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ all real spectral points of A belong to $\sigma_{ap}(A)$ (see e.g. [10, Corollary VI.6.2]).

Definition 2.1. For a self-adjoint operator A in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ a point $\lambda_0 \in \sigma(A)$ is called a spectral point of *positive (negative) type* of A if $\lambda_0 \in \sigma_{ap}(A)$ and for every sequence (x_n) in $\text{dom}(A)$ with $\|x_n\| = 1$, $n = 1, 2, \dots$, and $\|(A - \lambda_0)x_n\| \rightarrow 0$ for $n \rightarrow \infty$, we have

$$\liminf_{n \rightarrow \infty} [x_n, x_n]_{\mathcal{K}} > 0 \quad (\text{resp. } \limsup_{n \rightarrow \infty} [x_n, x_n]_{\mathcal{K}} < 0).$$

The point ∞ is said to be a point of *positive (negative) type* of the extended spectrum of A if A is unbounded and for every sequence (x_n) in $\text{dom}(A)$ with $\lim_{n \rightarrow \infty} \|x_n\| = 0$ and $\|Ax_n\| = 1$, $n = 1, 2, \dots$, we have

$$\liminf_{n \rightarrow \infty} [Ax_n, Ax_n]_{\mathcal{K}} > 0 \quad (\text{resp. } \limsup_{n \rightarrow \infty} [Ax_n, Ax_n]_{\mathcal{K}} < 0).$$

We denote the set of all points of $\tilde{\sigma}(A)$ of positive (negative) type by $\sigma_{++}(A)$ (resp. $\sigma_{--}(A)$). In the following proposition we collect some properties. For a proof we refer to [5].

Proposition 2.2.

- (i) *The sets $\sigma_{++}(A)$ and $\sigma_{--}(A)$ are contained in $\overline{\mathbb{R}}$.*
- (ii) *The non-real spectrum of A cannot accumulate to $\sigma_{++}(A) \cup \sigma_{--}(A)$.*
- (iii) *The sets $\sigma_{++}(A)$ and $\sigma_{--}(A)$ are relatively open in $\tilde{\sigma}(A)$.*
- (iv) *Let λ_0 be a point of $\sigma_{++}(A)$ ($\sigma_{--}(A)$, respectively). Then there exists an open neighbourhood \mathcal{U} in $\overline{\mathbb{C}}$ of λ_0 and a number $M > 0$ such that*

$$\|(A - \lambda)^{-1}\| \leq \frac{M}{|\text{Im } \lambda|} \quad \text{for all } \lambda \in \mathcal{U} \setminus \overline{\mathbb{R}}.$$

We shall say that an open subset Δ of $\overline{\mathbb{R}}$ is of *positive type (negative type)* with respect to A if

$$\Delta \cap \tilde{\sigma}(A) \subset \sigma_{++}(A) \quad (\text{resp. } \Delta \cap \tilde{\sigma}(A) \subset \sigma_{--}(A)).$$

An open set Δ of $\overline{\mathbb{R}}$ is called of *definite type* if Δ is of positive or negative type with respect to A . If we relate Definition 2.1 to nonnegative operators in Kreĭn spaces (cf. Section 2.1) we obtain from the properties of the spectral function of a nonnegative operator in a Kreĭn space, see, e.g., [3, 4, 39], and [5, Proposition 25] the following.

Proposition 2.3. *Let A be a nonnegative operator with $\rho(A) \neq \emptyset$ in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. Then $c(A) \subset \{0, \infty\}$ and*

$$\sigma(A) \cap (0, \infty) \subset \sigma_{++}(A) \subset \overline{\mathbb{R}} \setminus (-\infty, 0), \quad \sigma(A) \cap (-\infty, 0) \subset \sigma_{--}(A) \subset \overline{\mathbb{R}} \setminus (0, \infty).$$

In particular, we have

$$c(A) = \tilde{\sigma}(A) \setminus (\sigma_{++}(A) \cup \sigma_{--}(A)). \tag{2.2}$$

A generalization of the class of nonnegative operators in Kreĭn spaces is given by the class of definitizable operators. Recall, that a self-adjoint operator A in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ is called *definitizable* if $\rho(A) \neq \emptyset$ and if there exists a rational function $p \neq 0$ having poles only in $\rho(A)$ such that $[p(A)x, x]_{\mathcal{K}} \geq 0$ for all $x \in \mathcal{K}$. Such a function p is called *definitizing function* for A . Then the spectrum of A is real or its non-real part consists of a finite number of points. Inspired by Proposition 2.3 we introduce the set of *critical points* of a definitizable operator A via

$$c(A) := \tilde{\sigma}(A) \setminus (\sigma_{++}(A) \cup \sigma_{--}(A)). \quad (2.3)$$

It is known (cf. [39]) that $c(A)$ is contained in $\{t \in \mathbb{R} : p(t) = 0\} \cup \{\infty\}$.

For the definitizable operator A the spectral function $E(\Delta)$ can be introduced for every interval Δ such that the endpoints of Δ belong to intervals of definite type, see [39], [25]. We mention only that $E(\Delta)$ is defined and is a self-adjoint projection in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ for every such interval. Moreover,

$$(E(\Delta)\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}}) \text{ is a Hilbert space whenever } \Delta \subset \{t \in \mathbb{R} : p(t) > 0\}. \quad (2.4)$$

If a critical point α is the endpoint of two intervals (λ_1, α) and (α, λ_2) of definite type then the sequences $E([\lambda_1, t])$ and $E([t, \lambda_2])$ are monotone in (λ_1, α) and (α, λ_2) , resp. The point α is called a *regular critical point* of A , if the limits

$$\lim_{t \uparrow \alpha} E([\lambda_1, t]) \quad \text{and} \quad \lim_{t \downarrow \alpha} E([t, \lambda_2]) \quad (2.5)$$

exist in the strong operator topology. A critical point of A which is not regular is called *singular critical point* of A . The set of all singular critical points of A is denoted by $c_s(A)$.

In Subsection 4.2 we essentially use the following resolvent criterion of K. Veselić [46] for $\infty \notin c_s(A)$. We state a special case of this criterion as it has appeared in [23, Corollary 1.6].

Theorem 2.4. *Let A be a definitizable self-adjoint operator in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. Then:*

- (a) $\infty \notin c_s(A)$ if and only if there is $\eta_0 > 0$, such that the set of numbers

$$\int_{\eta_0}^{\eta} \operatorname{Re} [(A - iy)^{-1} f, f]_{\mathcal{K}} dy \quad (\eta \in (\eta_0, \infty))$$

is bounded for every $f \in \mathcal{K}$.

- (b) Let $\xi_0 \in \mathbb{R}$. Then $\xi_0 \notin c_s(A)$ and $\ker(A - \xi_0) = \ker(A - \xi_0)^2$ if and only if there is $\eta_0 > 0$, such that the set of numbers

$$\int_{\eta}^{\eta_0} \operatorname{Re} [(A - \xi_0 - iy)^{-1} f, f]_{\mathcal{K}} dy \quad (\eta \in (0, \eta_0))$$

is bounded for every $f \in \mathcal{K}$.

A characterization of definitizable operators via their sign-type spectrum together with some growth conditions for the resolvent is provided by the following theorem. Its proof follows from [26, Definition 4.4 and Theorem 4.7]).

Theorem 2.5. *Let A be a self-adjoint operator in the Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. Then A is definitizable if and only if the following holds.*

- (i) *The non-real spectrum $\sigma(A) \setminus \mathbb{R}$ consists of isolated points which are poles of the resolvent of A , and no point of $\overline{\mathbb{R}}$ is an accumulation point of the non-real spectrum $\sigma(A) \setminus \mathbb{R}$ of A .*

- (ii) There is an open neighbourhood \mathcal{U} of $\overline{\mathbb{R}}$ in $\overline{\mathbb{C}}$ and numbers $m \geq 1$, $M > 0$ with

$$\|(A - \lambda)^{-1}\| \leq M(|\lambda| + 1)^{2m-2} |\operatorname{Im} \lambda|^{-m} \quad \text{for all } \lambda \in \mathcal{U} \setminus \overline{\mathbb{R}}.$$

- (iii) Every point $\lambda \in \overline{\mathbb{R}}$ has an open connected neighbourhood I_λ in $\overline{\mathbb{R}}$ such that both components of $I_\lambda \setminus \{\lambda\}$ are of definite type with respect to A .

3. LOCALLY DEFINITIZABLE OPERATORS AND THEIR DIRECT SUM

3.1. Locally definitizable operators in Kreĭn spaces. In view of Theorem 2.5 it is natural to introduce a local version of definitizability which will play an important role in the following. The next notion is due to P. Jonas, see [24, 25], we mention also the overview in [45].

Definition 3.1. Let Ω be a domain in $\overline{\mathbb{C}}$ which is symmetric with respect to \mathbb{R} such that $\Omega \cap \overline{\mathbb{R}} \neq \emptyset$ and the intersections with the open upper and lower half-plane are simply connected. Let A be a self-adjoint operator in the Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. The operator A is called *definitizable over Ω* if the following holds.

- (i) The non-real spectrum in Ω , i.e. $\sigma(A) \cap (\Omega \setminus \overline{\mathbb{R}})$, consists of isolated points which are poles of the resolvent of A , and no point of $\Omega \cap \overline{\mathbb{R}}$ is an accumulation point of the non-real spectrum $\sigma(A) \setminus \mathbb{R}$ of A .
- (ii) For every closed subset Δ of $\Omega \cap \overline{\mathbb{R}}$ there exist an open neighborhood \mathcal{U} of Δ in $\overline{\mathbb{C}}$ and numbers $m \geq 1$, $M > 0$ such that

$$\|(A - \lambda)^{-1}\| \leq M(|\lambda| + 1)^{2m-2} |\operatorname{Im} \lambda|^{-m} \quad \text{for all } \lambda \in \mathcal{U} \setminus \overline{\mathbb{R}}.$$

- (iii) Every point $\lambda \in \Omega \cap \overline{\mathbb{R}}$ has an open connected neighborhood I_λ in $\overline{\mathbb{R}}$ such that both components of $I_\lambda \setminus \{\lambda\}$ are of definite type with respect to A .

Let A be definitizable over Ω . Similar as in (2.3) we call a point $t \in \Omega \cap \overline{\mathbb{R}}$ a *critical point* of the operator A if there is no open subset Δ of definite type with $t \in \Delta$. The set of critical points of A is denoted by $c(A)$. Similar as in Section 2.1 critical points admit a classification into singular and regular critical points: If for some $\lambda \in c(A) \setminus \{\infty\}$ the limits analogous to (2.5) exist, then λ is called a *regular critical point of A* . If ∞ is a critical point of A and the limits (2.5) exist in the strong operator topology for some $\lambda_1, \lambda_2 \in \mathbb{R} \setminus \{0\}$, then ∞ is called *regular critical point of A* . A critical point of A which is not regular is called *singular critical point of A* . The set of all singular critical points of A is denoted by $c_s(A)$.

Theorem 2.4 has a counterpart for locally definitizable operators: Let A be definitizable over a neighbourhood Ω of ∞ . Then A admits an orthogonal decomposition into two operators: a definitizable one with spectrum in $\overline{\Delta}$ and a self-adjoint one with spectrum outside Δ , where $\Delta(\subset \Omega)$ is a neighbourhood of ∞ , for details we refer to [26, Theorem 4.8]. Then the following theorem follows easily from this decomposition and Theorem 2.4.

Theorem 3.2. Let a self-adjoint operator A in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot])$ be locally definitizable over a neighborhood Ω of ∞ . Then $\infty \notin c_s(A)$ if and only if there is $\eta_0 > 0$, such that the set of numbers

$$\int_{\eta_0}^{\eta} \operatorname{Re} [(A - iy)^{-1} f, f]_{\mathcal{K}} dy \quad (\eta \in (\eta_0, \infty))$$

is bounded for every $f \in \mathcal{K}$.

Similarly, if $\xi_0 \in \mathbb{R}$ and A is locally definitizable over a neighborhood Ω of ξ_0 , then $\xi_0 \notin c_s(A)$ and $\ker(A - \xi_0) = \ker(A - \xi_0)^2$ if and only if there is $\eta_0 > 0$, such that the set of numbers

$$\int_{\eta}^{\eta_0} \operatorname{Re} [(A - \xi_0 - iy)^{-1} f, f]_{\mathcal{K}} dy \quad (\eta \in (0, \eta_0))$$

is bounded for every $f \in \mathcal{K}$.

Roughly speaking, the property of an operator to be definitizable or to be locally definitizable is stable under finite rank perturbations. This is made more precise in the following theorem which is taken from J. Behrndt [6, Theorem 2.2].

Theorem 3.3. *Let A_0 and A_1 be self-adjoint operators in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ with $\rho(A_0) \cap \rho(A_1) \neq \emptyset$ and assume that for some $\lambda_0 \in \rho(A_0) \cap \rho(A_1)$ the difference*

$$(A_0 - \lambda_0)^{-1} - (A_1 - \lambda_0)^{-1}$$

is a finite rank operator. Then A_0 is definitizable over Ω if and only if A_1 is definitizable over Ω .

Moreover, if A_0 is definitizable over Ω and $\delta \subset \Omega \cap \overline{\mathbb{R}}$ is an open interval with endpoint $\mu \in \Omega \cap \overline{\mathbb{R}}$ and the spectral points of A_0 in δ are only of positive type (negative type), then there exists an open interval δ' , $\delta' \subset \delta$, with endpoint μ such that the spectral points of A_1 in δ' are only of positive type (negative type, respectively).

Theorem 3.3 also holds for definitizable operators as the class of definitizable operators over $\overline{\mathbb{C}}$ coincides with the class of definitizable operators ([26, Theorem 4.7]). For definitizable operators this fact is already contained in [27].

3.2. Local definitizability of the direct sum of two operators. In this section we characterize the definitizability of an operator which is the direct sum of two definitizable operators. For this we provide the following definition.

Definition 3.4. We shall say that the sets S_1 and S_2 , $S_1, S_2 \subset \overline{\mathbb{R}}$, are separated by a finite number of points if there exists a finite ordered set $\{\alpha_j\}_{j=1}^N$, $N \in \mathbb{N}$,

$$-\infty = \alpha_0 < \alpha_1 \leq \cdots \leq \alpha_N < \alpha_{N+1} = +\infty,$$

such that one of the sets S_j , $j = 1, 2$, is a subset of $\bigcup_{k \text{ is even}} [\alpha_k, \alpha_{k+1}]$ and the other one is a subset of $\bigcup_{k \text{ is odd}} [\alpha_k, \alpha_{k+1}]$. Here we agree that 0 is even, $[\alpha_0, \alpha_1]$ stands for $(-\infty, \alpha_1] \cup \{\infty\}$ and $[\alpha_N, \alpha_{N+1}]$ for $[\alpha_N, \infty) \cup \{\infty\}$.

The following theorem can be considered as a refinement of [32, Theorem 3.6].

Theorem 3.5. *Consider two operators A and B where A is self-adjoint in the Kreĭn space $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}_+})$ and B in $(\mathcal{K}_-, [\cdot, \cdot]_{\mathcal{K}_-})$. Let the direct sum of the two Kreĭn spaces*

$$\mathcal{K} = \mathcal{K}_+ [+] \mathcal{K}_-$$

be endowed with the natural inner product

$$[f, g]_{\mathcal{K}} := [P_+ f, P_+ g]_{\mathcal{K}_+} + [P_- f, P_- g]_{\mathcal{K}_-} \quad (f, g \in \mathcal{K}), \quad (3.1)$$

where P_{\pm} are the orthogonal projections onto \mathcal{K}_{\pm} . Then the sum of the operators $A[+]B$ is self-adjoint in the direct sum of the Krein spaces \mathcal{K} with the natural inner product from (3.1). Set

$$S_+ := \sigma_{++}(A) \cup \sigma_{++}(B) \quad \text{and} \quad S_- := \sigma_{--}(A) \cup \sigma_{--}(B).$$

Then $A[+]B$ is definitizable if and only if the operators A and B are definitizable and S_+ and S_- are separated by a finite number of points.

Proof. The non-real-spectrum of $A[+]B$ coincides with the union of the non-real spectra of A and of B . Therefore, if $A[+]B$ is definitizable, then item (i) of Theorem 2.5 holds for A and for B . Conversely, if A and B are both definitizable, then (i) of Theorem 2.5 holds for $A[+]B$. Therefore, it is no restriction to assume that $A[+]B$, A , and B have real spectrum only.

If $A[+]B$ is definitizable, then by the definition of the inner product in $\mathcal{K} = \mathcal{K}_+[+] \mathcal{K}_-$ a definitizing function p for $A[+]B$ is also a definitizing function for A and for B . From (2.4) we deduce

$$\begin{aligned} \{t \in \mathbb{R} : p(t) > 0\} &\subset \sigma_{++}(A) \cup \rho(A), & \{t \in \mathbb{R} : p(t) < 0\} &\subset \sigma_{--}(A) \cup \rho(A), \\ \{t \in \mathbb{R} : p(t) > 0\} &\subset \sigma_{++}(B) \cup \rho(B), & \{t \in \mathbb{R} : p(t) < 0\} &\subset \sigma_{--}(B) \cup \rho(B), \end{aligned}$$

and, hence, the zeros of p are the points separating S_+ and S_- , cf. Definition 3.4.

It remains to prove the converse. Assume that S_+ and S_- are separated by the points $\{\alpha_0, \dots, \alpha_{N+1}\}$, cf. Definition 3.4, then we have

$$S_+ \cap S_- \subset \{\alpha_0, \dots, \alpha_{N+1}\}.$$

Note that S_+ and $c(A)$ may have an non-empty intersection (and the same applies to $S_+ \cap c(B)$, $S_- \cap c(A)$, and $S_- \cap c(B)$). Indeed, let $\lambda \in \sigma_{++}(B)$ (and, hence, $\lambda \in S_+$) such that λ is an isolated spectral point of A which belongs to $c(A)$. Then $\lambda \in S_+ \cap c(A)$ and, moreover as $\lambda \notin S_-$, we have in addition $\lambda \notin \{\alpha_0, \dots, \alpha_{N+1}\}$.

We define

$$\Lambda := \{\alpha_0, \dots, \alpha_{N+1}\} \cup c(A) \cup c(B)$$

and for $\lambda \in S_+ \setminus \Lambda$ the following statements are true.

- (i) $\lambda \in \sigma_{++}(A) \cup \sigma_{++}(B)$ (as $\lambda \in S_+$),
- (ii) $\lambda \notin \sigma_{--}(A) \cup \sigma_{--}(B)$ (as $\lambda \notin S_-$),
- (iii) $\lambda \notin c(A) \cup c(B)$ (as $\lambda \notin \Lambda$).

Thus, by (2.2) applied to both A and B , we obtain

$$\lambda \in \sigma_{++}(A) \cup \tilde{\rho}(A) \quad \text{and} \quad \lambda \in \sigma_{++}(B) \cup \tilde{\rho}(B).$$

This implies

$$\lambda \in \sigma_{++}(A[+]B)$$

and we obtain

$$S_+ \setminus \Lambda \subset \sigma_{++}(A[+]B) \tag{3.2}$$

and with similar arguments,

$$S_- \setminus \Lambda \subset \sigma_{--}(A[+]B). \tag{3.3}$$

From (2.2) we conclude

$$\begin{aligned} \tilde{\sigma}(A[+]B) &= \tilde{\sigma}(A) \cup \tilde{\sigma}(B) \\ &= \sigma_{++}(A) \cup c(A) \cup \sigma_{--}(A) \cup \sigma_{++}(B) \cup c(B) \cup \sigma_{--}(B) \\ &= S_+ \cup c(A) \cup c(B) \cup S_- \subset S_+ \cup S_- \cup \Lambda. \end{aligned} \tag{3.4}$$

Obviously, for the operator $A[+]B$ the statements (i) and (ii) from Theorem 2.5 are satisfied as A and B are definitizable operators. It remains to show (iii). Clearly, for $\lambda \in \overline{\mathbb{C}} \setminus \tilde{\sigma}(A[+]B)$ (iii) in Theorem 2.5 is satisfied. Let $\lambda \in \tilde{\sigma}(A[+]B)$. If $\lambda \in (S_+ \cup S_-) \setminus \Lambda$ we deduce from (3.2) and (3.3) that either $\lambda \in \sigma_{++}(A[+]B)$ or $\lambda \in \sigma_{--}(A[+]B)$. As the sets $\sigma_{++}(A[+]B)$ and $\sigma_{--}(A[+]B)$ are relatively open in $\tilde{\sigma}(A[+]B)$ (cf. Proposition 2.2), (iii) follows. By (3.4), it remains to consider $\lambda \in \Lambda$. For $\lambda \in \{\alpha_0, \dots, \alpha_{N+1}\}$ (iii) follows from (3.2) and (3.3). Therefore, consider $\lambda \in c(A) \cup c(B)$. It is sufficient to consider $\lambda \in c(A) \setminus \{\alpha_0, \dots, \alpha_{N+1}\}$. It follows from the definition of the points $\{\alpha_0, \dots, \alpha_{N+1}\}$ and the fact that $\lambda \notin \{\alpha_0, \dots, \alpha_{N+1}\}$ that there exists open connected neighbourhoods I_λ, J_λ in $\overline{\mathbb{R}}$ of λ with

$$(I_\lambda \setminus \{\lambda\}) \cap \tilde{\sigma}(A) \subset \sigma_{++}(A) \quad \text{and} \quad (J_\lambda \setminus \{\lambda\}) \cap \tilde{\sigma}(B) \subset \sigma_{++}(B)$$

or

$$(I_\lambda \setminus \{\lambda\}) \cap \tilde{\sigma}(A) \subset \sigma_{--}(A) \quad \text{and} \quad (J_\lambda \setminus \{\lambda\}) \cap \tilde{\sigma}(B) \subset \sigma_{--}(B).$$

This shows $(I_\lambda \cap J_\lambda \setminus \{\lambda\}) \cap \tilde{\sigma}(A[+]B)$ is a subset of $\sigma_{++}(A[+]B)$ or of $\sigma_{--}(A[+]B)$ and (iii) follows. \square

Corollary 3.6. *Let A_+ and A_- be self-adjoint and semibounded from below in the Kreĭn spaces $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}_+})$ and $(\mathcal{K}_-, [\cdot, \cdot]_{\mathcal{K}_-})$, respectively,*

$$[A_\pm f_\pm, f_\pm]_{\mathcal{K}_\pm} \geq \alpha_\pm [f_\pm, f_\pm]_{\mathcal{K}_\pm}, \quad f_\pm \in \text{dom}(A_\pm) \quad (3.5)$$

for some $\alpha_\pm \in \mathbb{R}$. Let $\rho(A_+) \neq \emptyset$, $\rho(A_-) \neq \emptyset$. Then their direct sum $A_+[+]A_-$ is definitizable over

$$\Omega := \overline{\mathbb{C}} \setminus [\min\{\alpha_+, \alpha_-\}, \max\{\alpha_+, \alpha_-\}] \quad (3.6)$$

in the direct sum of the Kreĭn spaces $\mathcal{K} = \mathcal{K}_+[+] \mathcal{K}_-$. In particular, $A_+[+]A_-$ is definitizable if and only if the sets S_+ and S_- from Theorem 3.5 are separated by a finite number of points.

This is fulfilled in the following special cases.

- (i) $\alpha_- = \alpha_+$.
- (ii) $\alpha_- < \alpha_+$ and either $\sigma(A_+) \cap (\alpha_-, \alpha_+)$ is finite or $\sigma(A_-) \cap (\alpha_-, \alpha_+)$ is finite.
- (iii) $\alpha_+ < \alpha_-$ and either $\sigma(A_+) \cap (\alpha_+, \alpha_-)$ is finite or $\sigma(A_-) \cap (\alpha_+, \alpha_-)$ is finite.

Proof. The assumptions on A_\pm imply that $A_+ - \alpha_+$ and $A_- - \alpha_-$ are nonnegative operators and, hence, A_\pm are definitizable operators. Then, with Proposition 2.3 we see that

$$(\alpha_\pm, \infty) \cap \sigma(A_\pm) \subset \sigma_{++}(A_\pm) \quad \text{and} \quad (-\infty, \alpha_\pm) \cap \sigma(A_\pm) \subset \sigma_{--}(A_\pm) \quad (3.7)$$

and properties (i)–(iii) from Definition 3.1 for the operator $A_+[+]A_-$ and Ω as in (3.6) are easily shown, cf. Proposition 2.2. Therefore, $A_+[+]A_-$ is definitizable over Ω .

The statements on the definitizability of the operator $A_+[+]A_-$ now follow directly from (3.7) and Theorem 3.5. \square

4. COUPLING OF DEFINITIZABLE OPERATORS IN KREĬN SPACES

4.1. Boundary triples and Weyl functions of symmetric operators. Starting from this section we will denote by A a closed densely defined symmetric operator in a Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. Let $\widehat{\rho}(A)$ denotes the set of points of regular type of A , see [1] and let \mathfrak{N}_z denote the defect subspace of the operator A

$$\mathfrak{N}_z := \mathcal{H} \ominus \text{ran}(A - \bar{z}) = \ker(A^+ - z), \quad z \in \widehat{\rho}(A).$$

In what follows we assume that the operator A admits a self-adjoint extension \tilde{A} in $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ with a nonempty resolvent set $\rho(\tilde{A})$. Then for all $z \in \rho(\tilde{A})$ we have

$$\text{dom}(A^+) = \text{dom}(\tilde{A}) + \mathfrak{N}_z \quad \text{direct sum in } \mathcal{H}. \quad (4.1)$$

This implies, in particular, that the dimension $\dim(\mathfrak{N}_z)$ is constant for all $z \in \rho(\tilde{A})$.

Definition 4.1. Let Γ_0 and Γ_1 be linear mappings from $\text{dom}(A^+)$ to \mathbb{C}^d such that

- (i) the mapping $\Gamma : f \rightarrow \{\Gamma_0 f, \Gamma_1 f\}$ from $\text{dom}(A^+)$ to \mathbb{C}^{2d} is surjective;
- (ii) the abstract Green's identity

$$[A^+ f, g]_{\mathcal{K}} - [f, A^+ g]_{\mathcal{K}} = (\Gamma_0 g)^*(\Gamma_1 f) - (\Gamma_1 g)^*(\Gamma_0 f) \quad (4.2)$$

holds for all $f, g \in \text{dom}(A^+)$.

Then the triplet $\Pi = \{\mathbb{C}^d, \Gamma_0, \Gamma_1\}$ is said to be a *boundary triple* for A^+ , see [21, 17, 18, Sect.3.1.4] for much more general setting.

It follows from (4.2) that the extensions A_0, A_1 of A defined as restrictions of A^+ to the domains

$$\text{dom}(A_0) := \ker(\Gamma_0) \quad \text{and} \quad \text{dom}(A_1) := \ker(\Gamma_1) \quad (4.3)$$

are self-adjoint extensions of A .

If A has a self-adjoint extension, lets say \tilde{A} , with $\rho(\tilde{A}) \neq \emptyset$, then the operator A^+ admits a boundary triple with $d = \dim \mathfrak{N}_z$, such that $A_0 = \tilde{A}$ and $d = \dim \mathfrak{N}_z$ ($z \in \rho(A_0)$). In this case for every $z \in \rho(A_0)$ the decomposition (4.1) holds with $\tilde{A} = A_0$ and the mapping $\Gamma_0|_{\mathfrak{N}_z}$ is invertible for every $z \in \rho(A_0)$. Therefore, the operator-function

$$\gamma(z) := (\Gamma_0|_{\mathfrak{N}_z})^{-1} \quad (4.4)$$

is well defined and takes values in $\mathcal{B}(\mathbb{C}^d, \mathfrak{N}_z)$. The operator-function $\gamma(z)$ is called the γ -field of A , associated with the boundary triple Π . Notice, that $\gamma(z)$ satisfies the equality

$$\gamma(z) = (A_0 - z)(A_0 - z)^{-1}\gamma(z_0) \quad (z, z_0 \in \rho(A_0)).$$

Definition 4.2. The matrix valued function $M : \rho(A_0) \rightarrow \mathbb{C}^{d \times d}$ is defined by the equality

$$M(z)\Gamma_0 f_z = \Gamma_1 f_z, \quad f_z \in \mathfrak{N}_z, z \in \rho(A_0). \quad (4.5)$$

The matrix valued function M is called the *Weyl function* of A corresponding to the boundary triplet $\Pi = \{\mathbb{C}^d, \Gamma_0, \Gamma_1\}$.

Clearly,

$$M(z) = \Gamma_1 \gamma(z), \quad z \in \rho(A_0), \quad (4.6)$$

and hence $M(z)$ is well defined and takes values in $\mathbb{C}^{d \times d}$. It follows from the identity that the Weyl function $M(\lambda)$ satisfies the identities

$$M(z) - M(w)^* = (z - \bar{w})\gamma(w)^+ \gamma(z), \quad z, w \in \rho(A_0). \quad (4.7)$$

With $w = \bar{z}$ the identity (4.7) yields that the Weyl function M satisfies the symmetry condition

$$M(\bar{z})^* = M(z) \quad \text{for all } z \in \rho(A_0). \quad (4.8)$$

The identity (4.7) was used in [35] as a definition of the Q -function. In the case when $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ is a Hilbert space it follows from (4.7) and (4.8) that M is a Nevanlinna function; cf. (1.2).

In what follows the function

$$\hat{f}(z) := [f, \gamma(\bar{z})]_{\mathcal{K}} \quad (f \in \mathcal{K}, z \in \rho(A_0))$$

is called the generalized Fourier transform of f associated with the boundary triplet $\{\mathbb{C}, \Gamma_0, \Gamma_1\}$. A motivation for this name is hidden in the fact, that the mapping $f \mapsto \hat{f}$ is a unitary mapping from \mathcal{K} to a reproducing kernel Kreĭn space with the kernel $\frac{M(z) - M(\bar{w})}{z - \bar{w}}$ (see [15] for the Hilbert space case).

Proposition 4.3. [35, 17, 18] *Let A_1 be the self-adjoint extension of A with the domain defined in (4.3) and let $d = 1$. For every $z \in \rho(A_0)$ the following equivalence hold:*

$$z \in \rho(A_1) \iff M(z) \neq 0$$

and the resolvent of A_1 can be found by the formula

$$(A_1 - z)^{-1} f = (A_0 - z)^{-1} f - \frac{\hat{f}(z)}{M(z)} \gamma(z)$$

for all $f \in \mathcal{H}$ and all $z \in \rho(A_0) \cap \rho(A_1)$.

4.2. Construction of the coupling of two self-adjoint operators in a Kreĭn space. In this section we consider two Kreĭn spaces $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}_+})$ and $(\mathcal{K}_-, [\cdot, \cdot]_{\mathcal{K}_-})$. Let their direct sum

$$\mathcal{K} = \mathcal{K}_+ [+] \mathcal{K}_-$$

be endowed with the natural inner product (3.1). Consider two closed symmetric densely defined operators A_+ and A_- with defect numbers $(1, 1)$ acting in the Kreĭn spaces $(\mathcal{K}_+, [\cdot, \cdot]_{\mathcal{K}_+})$ and $(\mathcal{K}_-, [\cdot, \cdot]_{\mathcal{K}_-})$. Let $(\mathbb{C}, \Gamma_0^\pm, \Gamma_1^\pm)$ be a boundary triple for A_\pm^\pm . Let M_\pm be the corresponding Weyl function and γ_{A_\pm} the γ -field. By $A_{\pm,0}$ we denote the self-adjoint extension of A_\pm which is defined on

$$\text{dom}(A_{\pm,0}) = \ker(\Gamma_0^\pm) \quad \text{by} \quad A_{\pm,0} = A_\pm^\pm|_{\ker(\Gamma_0^\pm)}.$$

and assume that $\rho(A_{+,0}) \cap \rho(A_{-,0}) \neq \emptyset$. Then the functions M_\pm are defined and holomorphic on $\rho(A_{\pm,0})$.

The following theorem is the indefinite version of a result from [19], (see also [15]).

Theorem 4.4. *Under the general assumptions of this subsection we have:*

(a) *The linear operator A defined as the restriction of $A_+^\pm [+] A_-^\pm$ to the domain*

$$\text{dom}(A) = \left\{ \begin{pmatrix} f_+ \\ f_- \end{pmatrix} : \begin{array}{l} \Gamma_0^+(f_+) = \Gamma_0^-(f_-) = 0, \\ \Gamma_1^+(f_+) + \Gamma_1^-(f_-) = 0, \end{array} f_\pm \in \text{dom}(A_\pm^\pm) \right\} \quad (4.9)$$

is closed, densely defined and symmetric with defect numbers $(1, 1)$ in the Kreĭn space \mathcal{K} .

(b) *The adjoint A^+ of A is the restriction of $A_+^\pm [+] A_-^\pm$ to the domain*

$$\text{dom}(A^+) = \left\{ \begin{pmatrix} f_+ \\ f_- \end{pmatrix} : \Gamma_0^+(f_+) - \Gamma_0^-(f_-) = 0, f_\pm \in \text{dom}(A_\pm^\pm) \right\}. \quad (4.10)$$

(c) A boundary triple $(\mathbb{C}, \Gamma_0, \Gamma_1)$ for A^+ is given by

$$\Gamma_0 f = \Gamma_0^+ f_+, \quad \Gamma_1 f = \Gamma_1^+ f_+ + \Gamma_1^- f_-, \quad f = \begin{pmatrix} f_+ \\ f_- \end{pmatrix} \in \text{dom}(A^+). \quad (4.11)$$

(d) The Weyl function $M(z)$ and the γ -field of A relative to the boundary triple $(\mathbb{C}, \Gamma_0, \Gamma_1)$ are given by

$$M(z) = M_+(z) + M_-(z), \quad \gamma(z) = \begin{pmatrix} \gamma_{A_+}(z) \\ \gamma_{A_-}(z) \end{pmatrix} \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (4.12)$$

(e) The self-adjoint extension A_1 of A such that $\text{dom}(A_1) = \ker(\Gamma_1)$ coincides with the restriction of $A_+^+[+]A_-^+$ to the domain

$$\text{dom}(A_1) = \left\{ \begin{pmatrix} f_+ \\ f_- \end{pmatrix} : \begin{array}{l} \Gamma_0^+(f_+) - \Gamma_0^-(f_-) = 0, \\ \Gamma_1^+(f_+) + \Gamma_1^-(f_-) = 0 \end{array} f_\pm \in \text{dom}(A_\pm^+) \right\}, \quad (4.13)$$

and is called a coupling of A_+ and A_- relative to the boundary triples $(\mathbb{C}, \Gamma_0^+, \Gamma_1^+)$ and $(\mathbb{C}, \Gamma_0^-, \Gamma_1^-)$.

(f) The self-adjoint extension A_0 of A coincides with the direct sum $A_{+,0}[+]A_{-,0}$ and $\rho(A_0) = \rho(A_{+,0}) \cap \rho(A_{-,0}) \neq \emptyset$.

(g) The resolvent set $\rho(A_1)$ is nonempty if and only if

$$M_+ + M_- \not\equiv 0.$$

For every $z \in \rho(A_1) \cap \rho(A_0)$ and $f = \begin{pmatrix} f_+ \\ f_- \end{pmatrix} \in \mathcal{K} = \mathcal{K}_+[+] \mathcal{K}_-$ the resolvent of A_1 is given by

$$(A_1 - z)^{-1} f = (A_0 - z)^{-1} f - \frac{\widehat{f}_{A_+}(z) + \widehat{f}_{A_-}(z)}{M_+(z) + M_-(z)} \gamma(z), \quad (4.14)$$

where

$$\widehat{f}_{A_+}(z) := [f_+, \gamma_{A_+}(\bar{z})]_{\mathcal{K}_+}, \quad \widehat{f}_{A_-}(z) := [f_-, \gamma_{A_-}(\bar{z})]_{\mathcal{K}_-}. \quad (4.15)$$

Proof. (a)–(c) Since $(\mathbb{C}, \Gamma_0^\pm, \Gamma_1^\pm)$ is a boundary triple for A_\pm^+ it follows from (4.2) that for all $f_\pm \in \text{dom}(A_\pm^+)$

$$\begin{aligned} [A_+^+ f_+, g_+]_{\mathcal{K}_+} - [f_+, A_+^+ g_+]_{\mathcal{K}_-} + [A_-^+ f_-, g_-]_{\mathcal{K}_-} - [f_-, A_-^+ g_-]_{\mathcal{K}_-} \\ = \overline{(\Gamma_0^+ g_+)} (\Gamma_1^+ f_+) - \overline{(\Gamma_1^+ g_+)} (\Gamma_0^+ f_+) \\ + \overline{(\Gamma_0^- g_-)} (\Gamma_1^- f_-) - \overline{(\Gamma_1^- g_-)} (\Gamma_0^- f_-). \end{aligned} \quad (4.16)$$

Denote by T the restriction of $A_+^+[+]A_-^+$ to the set of the right hand side of (4.10). If $f = \begin{pmatrix} f_+ \\ f_- \end{pmatrix}, g = \begin{pmatrix} g_+ \\ g_- \end{pmatrix} \in \text{dom}(T)$ then

$$\Gamma_0^+ f_+ = \Gamma_0^- f_- \quad \text{and} \quad \Gamma_0^+ g_+ = \Gamma_0^- g_-$$

and hence one obtains from (4.16)

$$[Tf, g]_{\mathcal{K}} - [f, Tg]_{\mathcal{K}} = \overline{\Gamma_0^+ g_+} (\Gamma_1^+ f_+ + \Gamma_1^- f_-) - \overline{(\Gamma_1^+ g_+ + \Gamma_1^- g_-)} \Gamma_0^+ f_+. \quad (4.17)$$

Now it follows from (4.17) that A is a closed, densely defined and symmetric operator in the Kreĭn space \mathcal{K} , $T = A^+$ and a boundary triple for A^+ can be chosen in the form (4.11).

(d) The formulas for M and γ are implied by (4.11), (4.4) and (4.5).

(e)&(f) As $(\mathbb{C}, \Gamma_0, \Gamma_1)$ is a boundary triple for A^+ , the extension A_1 with $\text{dom}(A_1) = \ker(\Gamma_1)$ is a restriction of $A_+^+[+]A_-^+$. The formula (4.13) for the domain follows from $A_1 \subset A^+$ (see (4.10)) and $\text{dom}(A_1) = \ker(\Gamma_1)$. The statement (f) is immediate from (4.10) and (4.11).

(g) The statement (g) is implied by (4.12) and Proposition 4.3. \square

Remark 4.5. The construction in Theorem 4.4 shows that the coupling of two self-adjoint operators $A_{+,0}$ and $A_{-,0}$ is not uniquely defined. Namely, let alongside with the boundary triple $\Pi^- = \{\mathbb{C}, \Gamma_0^-, \Gamma_1^-\}$ another boundary triple $\tilde{\Pi}^- = \{\mathbb{C}, \tilde{\Gamma}_0^-, \tilde{\Gamma}_1^-\}$ be defined by

$$\tilde{\Gamma}_0^- = c\Gamma_0^-, \quad \tilde{\Gamma}_1^- = \bar{c}^{-1}\Gamma_1^-$$

for some non-zero $c \in \mathbb{C}$, $c \neq 1$. Then the extension \tilde{A}_1 defined as the restriction of $A_+^+[+]A_-^+$ to the domain

$$\text{dom}(\tilde{A}_1) = \left\{ \begin{pmatrix} f_+ \\ f_- \end{pmatrix} : \begin{array}{l} \Gamma_0^+(f_+) - c\Gamma_0^-(f_-) = 0, \\ \Gamma_1^+(f_+) + \bar{c}^{-1}\Gamma_1^-(f_-) = 0, \end{array} f_\pm \in \text{dom}(A_\pm^+) \right\}$$

is also a coupling of A_- and A_+ with $\tilde{A}_1 \neq A_1$.

However, when the boundary triples $\{\mathbb{C}, \Gamma_0^\pm, \Gamma_1^\pm\}$ are fixed then the coupling A_1 of the operators A_\pm is uniquely defined by the formula (4.13) and is called the *coupling* of the operators $A_{\pm,0}$ relative to the boundary triples $\{\mathbb{C}, \Gamma_0^\pm, \Gamma_1^\pm\}$.

Let us suppose that the operators $A_{\pm,0}$ are semibounded from below, that is there exists $\alpha_\pm \in \mathbb{R}$ such that (3.5) holds. Then the results of Section 3.2 allow to show that the coupling A_1 of the operators $A_{+,0}$ and $A_{-,0}$ is at least locally definitizable in a neighborhood of ∞ . In the next theorem sufficient conditions for regularity of the critical point ∞ are given.

Theorem 4.6. *Under the general assumptions of this subsection assume that the operators $A_{\pm,0}$, the γ -fields γ_\pm and the Weyl functions M_\pm satisfy the following assumptions:*

(A1) *The operators $A_{\pm,0}$ are semibounded from below, $\rho(A_{\pm,0}) \neq \emptyset$, and*

$$\infty \notin c_s(A_{\pm,0}).$$

(A2) *($w(z) := |M_+(z) + M_-(z)| \not\equiv 0$ on $\rho(A_{+,0}) \cap \rho(A_{-,0})$)*

(A3) *There is $y_1 > 0$, such that for all $f_{A_\pm} \in \mathcal{K}_\pm$*

$$\int_{y_1}^\infty \frac{|\widehat{f}_{A_\pm}(iy)|^2}{w(iy)} dy < \infty, \quad \int_{y_1}^\infty \frac{|\widehat{f}_{A_\pm}(-iy)|^2}{w(iy)} dy < \infty. \quad (4.18)$$

where the generalized Fourier transforms \widehat{f}_{A_+} and \widehat{f}_{A_-} are defined by (4.15).

Then the coupling A_1 of the operators $A_{+,0}$ and $A_{-,0}$ is definitizable over Ω , where Ω is as in (3.6). Moreover, we have

$$\infty \notin c_s(A_1).$$

Proof. By Corollary 3.6 the operator $A_0 = A_{+,0}[+]A_{-,0}$ is definitizable over Ω . In view of Theorem 4.4 the assumption (A2) yields $\rho(A_1) \neq \emptyset$. Since the operator A_1 is a two-dimensional perturbation of A_0 , by Theorem 3.3, the operator A_1 is also definitizable over Ω .

Clearly, $\infty \notin c_s(A_0)$ and it follows from Theorem 3.2 that there is $y_2 > y_1 > 0$, such that

$$\int_{y_2}^{\infty} |\operatorname{Re} [(A_0 - iy)^{-1} f, f]_{\mathcal{K}}| dy < \infty \quad \text{for all } f \in \mathcal{K}.$$

Let us set

$$\mathcal{A}(f, iy) := \frac{(\widehat{f}_{A_+}(iy) + \widehat{f}_{A_-}(iy))\overline{(\widehat{f}_{A_+}(iy) + \widehat{f}_{A_-}(iy))}}{M_+(iy) + M_-(iy)}. \quad (4.19)$$

We show

$$\int_{y_2}^{\infty} |\mathcal{A}(f, iy)| dy < \infty \quad \text{for all } f \in \mathcal{K}.$$

It follows from (A3) that for every $f_{A_{\pm}} \in \mathcal{K}_{\pm}$

$$\begin{aligned} & \int_{y_2}^{\infty} \left| \widehat{f}_{A_{\pm}}(iy) \widehat{f}_{A_{\pm}}(-iy) \right| \frac{dy}{w(iy)} \\ & \leq \left(\int_{y_2}^{\infty} \left| \widehat{f}_{A_{\pm}}(iy) \right|^2 \frac{dy}{w(iy)} \right)^{1/2} \left(\int_{y_2}^{\infty} \left| \widehat{f}_{A_{\pm}}(-iy) \right|^2 \frac{dy}{w(iy)} \right)^{1/2} < \infty. \end{aligned} \quad (4.20)$$

Similarly, one obtains for all $f_{A_{\pm}} \in \mathcal{K}_{\pm}$

$$\int_{y_2}^{\infty} \left| \widehat{f}_{A_+}(iy) \widehat{f}_{A_-}(-iy) \right| \frac{dy}{w(iy)} < \infty. \quad (4.21)$$

Combining (4.20) and (4.21) one obtains from (4.19) for all $f \in \mathcal{K}$

$$\int_{y_2}^{\infty} |\mathcal{A}(f, iy)| dy = \int_{y_2}^{\infty} \left| \frac{(\widehat{f}_{A_+}(iy) + \widehat{f}_{A_-}(iy))\overline{(\widehat{f}_{A_+}(-iy) + \widehat{f}_{A_-}(-iy))}}{M_+(iy) + M_-(iy)} \right| dy < \infty.$$

Now the statement $\infty \notin c_s(A_1)$ is implied by Theorem 2.4 and (4.14). \square

Theorem 4.7. *Under the assumptions of this subsection we assume that the operators $A_{\pm,0}$, the γ -fields γ_{\pm} and the Weyl functions M_{\pm} satisfy the following assumptions:*

(A1') *The operators $A_{\pm,0}$ are semibounded from below, $\rho(A_{\pm,0}) \neq \emptyset$, one of the conditions (i), (ii) or (iii) of Corollary 3.6 holds, and $\alpha := \min\{\alpha_-, \alpha_+\}$ satisfies*

$$\alpha \notin c_s(A_{\pm,0}).$$

(A2') *$(w(z) := |M_+(z) + M_-(z)| \not\equiv 0$ on $\rho(A_{+,0}) \cap \rho(A_{-,0})$.*

(A3') *There is $y_1 > 0$, such that for all $f_{A_{\pm}} \in \mathcal{K}_{\pm}$*

$$\int_0^{y_1} \frac{|\widehat{f}_{A_{\pm}}(\alpha + iy)|^2}{w(\alpha + iy)} dy < \infty, \quad \int_0^{y_1} \frac{|\widehat{f}_{A_{\pm}}(\alpha - iy)|^2}{w(\alpha - iy)} dy < \infty.$$

Then the coupling A_1 of the operators $A_{+,0}$ and $A_{-,0}$ is a definitizable operator and

$$\alpha \notin c_s(A).$$

Proof. In view of Corollary 3.6 the operator $A_0 := A_{+,0}[\dot{+}]A_{-,0}$ is definitizable. By Theorem 4.4 the assumption (A2') implies $\rho(A_!) \neq \emptyset$. Then by [27] the operator A_1 is also definitizable.

By the assumption (A1') $\alpha \notin c_s(A_{\pm,0})$ and hence $\alpha \notin c_s(A_0)$. Since by Theorem 2.4 there is $y_2 \in (0, y_1)$, such that

$$\int_0^{y_2} |\operatorname{Re} [(A_0 - \alpha - iy)^{-1} f, f]_{\mathcal{K}}| dy < \infty \quad \text{for all } f \in \mathcal{K}.$$

it remains to show that

$$\int_0^{y_2} |\mathcal{A}(f, \alpha + iy)| dy < \infty \quad \text{for all } f \in \mathcal{K},$$

where \mathcal{A} is defined as in (4.19). The proof of this inequality is similar to that in Theorem 4.6 and is based on the assumption (A3'). \square

5. APPLICATION TO STURM-LIOUVILLE OPERATORS WITH INDEFINITE WEIGHTS

Consider the differential expression

$$\ell(f)(t) := \frac{\operatorname{sgn} t}{w(t)} \left(-\frac{d}{dt} \left(\frac{df}{r(t)dt} \right) + q(t)f(t) \right) \quad \text{for a.a. } t \in \mathbb{R}, \quad (5.1)$$

where the coefficients r, q and w are real functions on \mathbb{R} satisfying the conditions:

- (C1) $r, q, w \in L^1_{\text{loc}}(\mathbb{R})$ and $r, w > 0$ a.e. on \mathbb{R} ,
- (C2) the expression ℓ is in the limit point case at $-\infty$ and at $+\infty$.

Let $\mathcal{H}_\pm = L_w^2(\mathbb{R}_\pm)$ be the standard weighted L^2 -space with the positive definite inner product

$$(f, g)_\pm = \int_{\mathbb{R}_\pm} f(t)\overline{g(t)}w(t)dt \quad (f, g \in L_w^2(\mathbb{R}_\pm)).$$

Consider minimal differential operators B_\pm generated by $\pm\ell$ in $L_{w_\pm}^2(\mathbb{R}_\pm)$, here w_\pm denotes the restriction of w to \mathbb{R}_\pm . Since we assume that ℓ is in the limit point case at $\pm\infty$, the operator B_\pm is a densely defined symmetric operator with defect numbers $(1, 1)$ in the Hilbert space $L_{w_\pm}^2(\mathbb{R}_\pm)$ and

$$\begin{aligned} \operatorname{dom}(B_\pm^*) &= \{f \in L_{w_\pm}^2(\mathbb{R}_\pm) : f, (r^{-1}f') \in AC_{\text{loc}}[0, \pm\infty), \ell(f) \in L_{w_\pm}^2(\mathbb{R}_\pm)\}, \\ \operatorname{dom}(B_\pm) &= \{f \in \operatorname{dom}(B_\pm^*) : f(0) = f'(0) = 0\}, \\ B_\pm f &:= \pm\ell(f), \quad f \in \operatorname{dom}(B_\pm). \end{aligned} \quad (5.2)$$

In addition to (C1), (C2) assume that

- (C3) B_+ and B_- are semibounded from below in $L_{w_+}^2(\mathbb{R}_+)$ and $L_{w_-}^2(\mathbb{R}_-)$, respectively.

Let $z \in \mathbb{C} \setminus \mathbb{R}$ and denote by $\vartheta(\cdot, z)$ and $\varphi(\cdot, z)$ the unique solutions of the equation

$$-(r^{-1}f')' + qf = zwf$$

satisfying the boundary conditions

$$\varphi(0, z) = 1, \quad (r^{-1}\varphi')(0, z) = 0 \quad \text{and} \quad \vartheta(0, z) = 0, \quad (r^{-1}\vartheta')(0, z) = 1,$$

respectively. Since we assume that $\pm\ell$ are in the limit point case at $\pm\infty$, for each $z \in \mathbb{C} \setminus \mathbb{R}$ there is a unique solution

$$\psi_\pm(t, z) = \varphi(t, z) \pm m_\pm(z)\vartheta(t, z), \quad t \in \mathbb{R}_\pm, \quad (5.3)$$

of the restriction of $\pm\ell(f) = zf$ to \mathbb{R}_\pm which belongs to $L_{w_\pm}^2(\mathbb{R}_\pm)$. Relation (5.3) defines the function $m_\pm : \mathbb{C} \setminus \mathbb{R} \rightarrow \mathbb{C}$ uniquely. The function m_\pm is called the *Dirichlet m-coefficient* of the restriction of the expression $\pm\ell$ to \mathbb{R}_\pm .

A boundary triple for B_\pm^* is $(\mathbb{C}, \Gamma_0^\pm, \Gamma_1^\pm)$, where

$$\Gamma_0^\pm f := f(0\pm), \quad \Gamma_1^\pm(f) = \pm(r^{-1}f')(0\pm), \quad f \in \operatorname{dom}(B_\pm^*). \quad (5.4)$$

It follows from (4.6) and (5.4) that the Dirichlet m -coefficient m_{\pm} defined by (5.3) coincides with the Weyl function of the operator B_{\pm} in (5.2) relative to the boundary triple in (5.4).

It is natural to consider the expression ℓ in the Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$, where $\mathcal{K} = L_w^2(\mathbb{R})$ is the standard weighted L^2 -space endowed with the indefinite inner product

$$[f, g]_{\mathcal{K}} = (Jf, g)_{L_w^2(\mathbb{R})} = \int_{\mathbb{R}} \operatorname{sgn} t f(t) \overline{g(t)} dt, \quad f, g \in L_w^2(\mathbb{R}),$$

and the operator

$$(Jf)(t) = (\operatorname{sgn} t) f(t), \quad f \in L_w^2(\mathbb{R}),$$

is a fundamental symmetry on $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. Set

$$\mathcal{K}_{\pm} = \{f \in L_w^2(\mathbb{R}) : f = 0 \text{ a.e. on } \mathbb{R}_{\mp}\}.$$

Then $\mathcal{K} = \mathcal{K}_+[\dot{+}]\mathcal{K}_-$ is the fundamental decomposition corresponding to J .

Let the operators $A_{\pm} := \pm B_{\pm}$ be considered as semibounded symmetric operators in Kreĭn spaces $(L_{w_{\pm}}^2(\mathbb{R}_{\pm}), \pm(\cdot, \cdot)_{L_{w_{\pm}}^2(\mathbb{R}_{\pm})})$. Then the triples (5.4) are boundary triples for A_{\pm}^+ . The corresponding Weyl functions of the operators A_+ and A_- take the form

$$M_+(z) = m_+(z), \quad M_-(z) = m_-(-z).$$

Consider a symmetric operator A in the Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$ determined by the conditions (4.9). Then the domain of the adjoint operator A^+ is characterized by the boundary condition (4.10), which in view of (5.4) takes the form

$$f(0+) = f(0-).$$

Consider the coupling A_1 of A_+ and A_- relative to the boundary triples (5.4). A_1 is characterized by the boundary conditions (4.13), which now can be rewritten as

$$f(0+) = f(0-), \quad (r^{-1}f')(0+) = (r^{-1}f')(0-).$$

Therefore, the operator A_1 is associated with the expression in (5.1) in the Hilbert space $L_w^2(\mathbb{R})$; that is $A_1 f = \ell(f)$ for all

$$f \in \operatorname{dom}(A_1) = \{f \in L_w^2(\mathbb{R}) : f, r^{-1}f' \in AC_{\operatorname{loc}}(\mathbb{R}), \ell(f) \in L_w^2(\mathbb{R})\}.$$

Notice, that the assumptions (A1) of Theorem 4.6 is satisfied in view of (C3) and the assumptions (A2) is satisfied since if $m_+(z) + m_-(-z) \equiv 0$ then $m_+(z) = -m_-(-z)$ is holomorphic on the half-line $(-\beta_-, \infty)$, what is impossible for the m -coefficient of the Sturm-Liouville operator. These considerations and Theorem 4.6 justify the following

Proposition 5.1. *Let the differential operation ℓ satisfy (C1), (C2) and let the minimal differential operators B_{\pm} generated by $\pm \ell$ in $L_w^2(\mathbb{R}_{\pm})$ satisfy (C3) and let m_{\pm} be the Dirichlet m -functions of B_{\pm} . Then the coupling A_1 of A_+ and A_- is locally definitizable in the Kreĭn space $(\mathcal{K}, [\cdot, \cdot]_{\mathcal{K}})$. If, in addition, m_+ and m_- satisfy the condition (4.18), then $\infty \notin c_s(A_1)$.*

REFERENCES

- [1] N.I. Akhiezer and I.M. Glazman, *Theory of Linear Operators in Hilbert Space*, two volumes bound as one, Dover Publications, 1993.
- [2] S. Albeverio, U. Günther and S. Kuzhel, J-self-adjoint operators with C-symmetries: extension theory approach, *J. Phys. A: Math. Theor.* 42 (2009), 105205.
- [3] T. Ando, *Linear Operators in Krein Spaces*, Lecture Notes, Hokkaido University, Sapporo, 1979.
- [4] T.Ya. Azizov and I.S. Iokhvidov, *Linear operators in spaces with an indefinite metric*, John Wiley & Sons, 1990.
- [5] T. Ya. Azizov, P. Jonas and C. Trunk, Spectral points of type π_+ and π_- of selfadjoint operators in Krein spaces, *J. Funct. Anal.* 226 (2005), 114-137.
- [6] J. Behrndt, Finite rank perturbations of locally definitizable self-adjoint operators in Krein spaces, *J. Operator Theory* 58 (2007), 415–440.
- [7] J. Behrndt, On the spectral theory of singular indefinite Sturm-Liouville operators, *J. Math. Anal. Appl.* 334 (2007) 1439–1449.
- [8] C. M. Bender, D. C. Brody and H. F. Jones, Complex extension of quantum mechanics, *Phys. Rev. Lett.* 89 (2002), 270401.
- [9] C. M. Bender, D. C. Brody and H. F. Jones, Must a Hamiltonian be Hermitian? *Am. J. Phys.* 71 (2003) 1095–1102.
- [10] J. Bognar, *Indefinite inner product spaces*, Springer, 1974.
- [11] V.M. Bruk, On a class of problems with the spectral parameter in the boundary conditions, *Mat. Sb.* 100 (1976), 210–216.
- [12] E. Caliceti, S. Graffi, M. Hitrik and J. Sjöstrand, Quadratic \mathcal{PT} -symmetric operators with real spectrum and similarity to self-adjoint operators, *J. Phys. A: Math. Theor.* 45 (2012), 444007.
- [13] E. Caliceti, S. Graffi and J. Sjöstrand, Spectra of \mathcal{PT} -symmetric operators and perturbation theory, *J. Phys. A: Math. Gen.* 38 (2005), 185–193.
- [14] J.W. Calkin, Abstract symmetric boundary conditions, *Trans. Am. Math. Soc.* 45 (1939), 369-442.
- [15] B. Ćurgus and V. Derkach, Partially fundamentally reducible operators in Krein spaces, *Integral Equations Operator Theory* 82 (2015) 469-518.
- [16] B. Ćurgus and H. Langer, A Krein space approach to symmetric ordinary differential operators with an indefinite weight function. *J. Differential Equations* 79 (1989) 31–61.
- [17] V.A. Derkach and M.M. Malamud, Generalized resolvents and the boundary value problems for hermitian operators with gaps, *J. Funct. Anal.* 95 (1991) 1–95.
- [18] V.A. Derkach, On Weyl function and generalized resolvents of a Hermitian operator in a Krein space, *Integral Equations Operator Theory* 23 (1995) 387–415.
- [19] V.A. Derkach, S. Hassi, M.M. Malamud and H.S.V. de Snoo, Generalized resolvents of symmetric operators and admissibility, *Methods Funct. Anal. Topol.* 6 (2000), 24–55.
- [20] W.F. Donoghue, *Monotone matrix functions and analytic continuation*. Die Grundlehren der mathematischen Wissenschaften, Band 207, Springer-Verlag, 1974.
- [21] V.I. Gorbachuk and M.L. Gorbachuk, *Boundary value problems for operator differential equations*, Kluwer Academic Publishers Group, 1991.
- [22] U. Günther, F. Stefani and M. Znojil, MHD α^2 -dynamo, squire equation and \mathcal{PT} -symmetric interpolation between square well and harmonic oscillator, *J. Math. Phys.* 46 (2005), 063504.
- [23] P. Jonas, Regularity criteria for critical points of definitizable operators. 179–195, *Operator Theory: Advances and Applications* 14, Birkhäuser, 1984.
- [24] P. Jonas, On a class of unitary operators in Krein space, *Oper. Theory Adv. Appl.* 17 (1986) 151–172.
- [25] P. Jonas, On a class of self-adjoint operators in Krein space and their compact perturbations, *Integral Equations Operator Theory* 11 (1988) 351–384.
- [26] P. Jonas, On Locally Definite Operators in Krein Spaces, in “Spectral Theory and its Applications,” Ion Colojoară Anniversary Volume, Theta, Bucharest, 2003, 95–127.
- [27] P. Jonas and H. Langer, Compact perturbations of definitizable operators, *J. Operator Theory* 2 (1979), 63–77.
- [28] I.S. Kac and M.G. Krein, R -functions–analytic functions mapping the upper halfplane into itself, *Amer. Math. Soc. Transl. Ser.* 103 (1974), 1–18.

- [29] I.M. Karabash, *J*-self-adjoint ordinary differential operators similar to self-adjoint operators. *Methods Funct. Anal. Topology* 6 (2000) 22–49.
- [30] I.M. Karabash, A. Kostenko, M.M. Malamud, The similarity problem for *J*-nonnegative Sturm-Liouville operators, *J. Differential Equations* 246 (2009) 964–997.
- [31] I.M. Karabash, M.M. Malamud, Indefinite Sturm-Liouville operators with finite zone potentials, *Operators and Matrices* 1 (2007) 301–368.
- [32] I. Karabash and C. Trunk, Spectral properties of singular Sturm-Liouville operators, *Proc. R. Soc. Edinb. Sect. A* 139 (2009), 483–503.
- [33] A.N. Kochubei, On extinctions of symmetric operators and symmetric binary relations, *Matem. Zametki* 17 (1975), 41–48.
- [34] A. Kostenko, The similarity problem for indefinite Sturm-Liouville operators and the HELP inequality, *Adv. Math.* 246 (2013), 368–413.
- [35] M.G. Krein and H. Langer, Über einige Fortsetzungsprobleme, die eng mit der Theorie hermitescher Operatoren in Raume Π_κ zusammenhängen, Teil I: Einige Funktionenklassen und ihre Darstellungen, *Math. Nachr.* 77 (1977), 187–236.
- [36] P. Lancaster, A.S. Markus and V.I. Matsaev, Definitizable operators and quasihyperbolic operator polynomials, *J. Funct. Anal.* 131 (1995), 1–28.
- [37] H. Langer, Spektraltheorie linearer Operatoren in *J*-Räumen und einige Anwendungen auf die Schar $L(\lambda) = \lambda^2 I + \lambda B + C$, *Habilitationsschrift*, Technische Universität Dresden (1965).
- [38] H. Langer, Verallgemeinerte Resolventen eines *J*-nichtnegativen Operators mit endlichem Defekt, *J. Functional Analysis* 8 (1971), 287–320.
- [39] H. Langer, Spectral functions of definitizable operators in Krein spaces. Functional analysis (Dubrovnik, 1981) 1–46, *Lecture Notes in Math.*, 948, Springer, 1982.
- [40] H. Langer, A.S. Markus and V.I. Matsaev, Locally definite operators in indefinite inner product spaces. *Math. Ann.* 308 (1997), 405–424.
- [41] H. Langer and C. Tretter, A Krein space approach to \mathcal{PT} -symmetry, *Czech. J. Phys.* 54 (2004), 1113–1120.
- [42] A.V. Straus, Extensions and generalized resolvents of a symmetric operator which is not densely defined, *Izv. Akad. Nauk. SSSR, Ser. Mat.* 34 (1970), 175–202 (Russian) (English translation: *Math. USSR-Izvestija*, 4 (1970), 179–208).
- [43] T. Tanaka, \mathcal{PT} -symmetric quantum theory defined in a Krein space, *J. Phys. A: Math. Gen.* 39 (2006), L369–L376.
- [44] T. Tanaka, General aspects of \mathcal{PT} -symmetric and *P*-self-adjoint quantum theory in a Krein space, *J. Phys. A: Math. Gen.* 39 (2006), 14175–14203.
- [45] C. Trunk, Locally definitizable operators: The local structure of the spectrum. in: *Operator Theory*, D. Alpay (Ed.), 241–259. Springer 2015.
- [46] K. Veselić, On spectral properties of a class of *J*-self-adjoint operators. I, II. *Glasnik Mat. Ser. III* 7 (1972) 229–248; ibid. 7 (1972) 249–254.

DEPARTMENT OF MATHEMATICS, DRAGOMANOV NATIONAL PEDAGOGICAL UNIVERSITY, KIEV,
AND DEPARTMENT OF MATHEMATICS, VASYL STUS DONETSK NATIONAL UNIVERSITY, VINNYTSYA,
UKRAINE

E-mail address: derkach.v@gmail.com

INSTITUT FÜR MATHEMATIK, TECHNISCHE UNIVERSITÄT ILMENAU, POSTFACH 100565, D-98684
ILMENAU, GERMANY

E-mail address: carsten.trunk@tu-ilmenau.de